

Photoinduced insulating of layered organic metal driven by strong electric field of 1.5-cycle, 7 fs infrared pulse

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Abstract: Photoinduced metal to insulator (M-I) change were realized by strong electric field (10 MV/cm) of 1.5-cycle (7 fs) near infrared pulse. Threshold for the M-I change decrease near the M-I transition temperature. The M-I change is driven by coherent charge oscillation with a period of 20 fs.

OCIS codes: (320.7120)Ultrafast Phenomena; (320.7130) Ultrafast processes in condensed matter, including semiconductors

1. Introduction

Photoinduced insulator to metal transition (PIMT) or equivalently the melting of the charge ordering (CO) insulators attract much attentions, because they can trigger various ultrafast changes in conducting and/ or magnetic properties. Such ultrafast melting of CO is driven by the immediate changes of the band filling and /or the electron kinetic energy. On the other hand, there are few studies on the reverse process, i.e., the metal to insulator (M-I) change. Considering the thermodynamic nature of the CO, i.e., the entropy for the CO is smaller than that for metal, photoinduced formation of the CO has been considered to be difficult. However, an instantaneous strong electric field >10 MV/cm which is realized by the few optical-cycle fs laser enables us to modulate the electronic structures directly before the increase of electron temperature.

The layered organic conductor α -(ET)₂I₃ (ET; bis[ethylenedithio]-tetrathiafulvalene) exhibiting ferroelectric CO insulator(Fig. 1(a))-metal(Fig. 1(b)) transition (transition temperature $T_{CO}=135$ K) has been the key material for investigating PIMT[1, 2]. In this study, we have demonstrated the photoinduced M-I change as shown schematic illustration in Figs. 1(a)(b) by using the 7 fs, 1.5-cycle near infrared pulse. Ultrafast reflectivity change showing the M-I change were discussed in terms of the "dynamical localization" [3-5].

2. Experiment

Super broadband infrared spectrum covering 1.2-2.3 μm was obtained by focusing the idler pulse (1.7 μm , carrier-envelope phase is self-stabilized) from the optical parametric amplifier to the hollow-fiber set in Kr filled chamber. Pulse compression was performed using both active mirror and chirped mirror techniques. Pulse width evaluated from the second harmonic generation autocorrelation is 7 fs which corresponds to 1.5-optical cycle. Instantaneous electric field on the sample surface (excitation diameter ~ 160 μm) reaches to > 10 MV/cm.

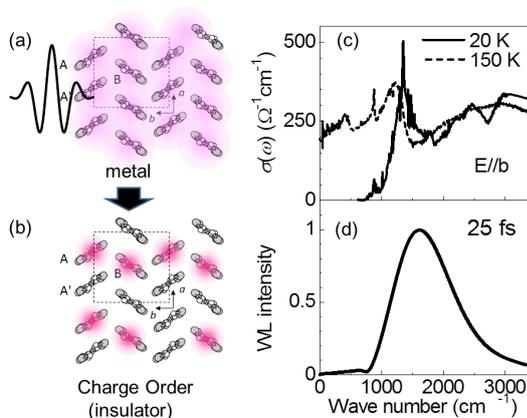


Fig. 1. (a), (b)Schematic illustrations of photoinduced change from metal(a) to charge ordering (CO) insulator(b). (c)Steady state optical conductivity. (d) Wavelet analysis of 20 fs oscillation in time profile shown in Fig. 3(See text).

3. Results and discussions

Fig. 2 shows the steady state reflectivity (R) spectrum (Fig. 2(a)), spectral differences between different temperatures (Fig. 2(b)), and transient reflectivity $\Delta R/R$ spectra (Fig. 2(c)), respectively. An increase of R (positive $\Delta R/R$) at 0.6-0.75 eV was observed at 35 fs after the excitation by 7 fs pulse (closed circles and gray shade in Fig. 2(c)), although R decreases in 0.6-0.65 eV at 300 fs, indicating the negative $\Delta R/R$. The spectral shape of $\Delta R/R$ at 30 fs is analogous to the temperature differential spectrum $[R(40\text{K})-R(140\text{K})]/R(140\text{K})$ in Fig 2(b), reflecting the decrease of temperature or the M-I change. On the other hand, $\Delta R/R$ spectrum at 300 fs is analogous to the $[R(250\text{K})-R(140\text{K})]/R(140\text{K})$, indicating the increase of temperatures. Fig. 3(a) shows the excitation intensity (I_{ex}) dependence of $\Delta R/R$ (138 K, 35 fs). Increase of R is detectable only for $I_{\text{ex}} > 0.5$ mJ/cm^2 (6.3 MV/cm), showing the threshold-like behaviour. The negative $\Delta R/R$

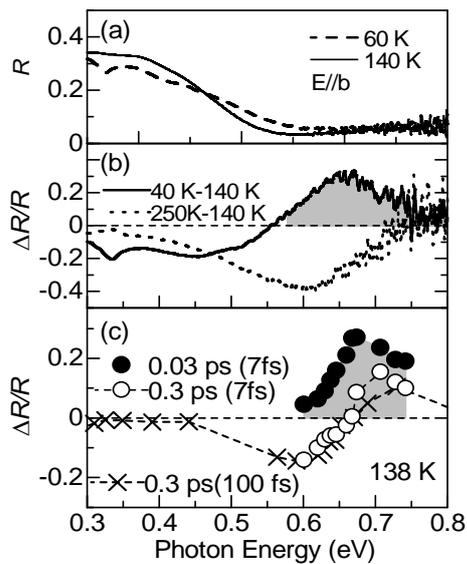


Fig. 2. (a) Steady state reflectivity spectra (E/b) of α -(ET)₂I₃ at 60 K (CO) and 140 K (metal). (b) Spectral difference $R(40\text{K})-R(140\text{K})/R(140\text{K})$ (solid curve) and $[R(250\text{K})-R(140\text{K})]/R(250\text{K})$ (dashed curve). (c) Transient reflectivity ($\Delta R/R$) spectra at 0.03 (closed) and 0.3 ps (open) after excitation by 7 fs pulse. Dashed curve and crosses show the $\Delta R/R$ at 0.3 ps after 100 fs excitation.

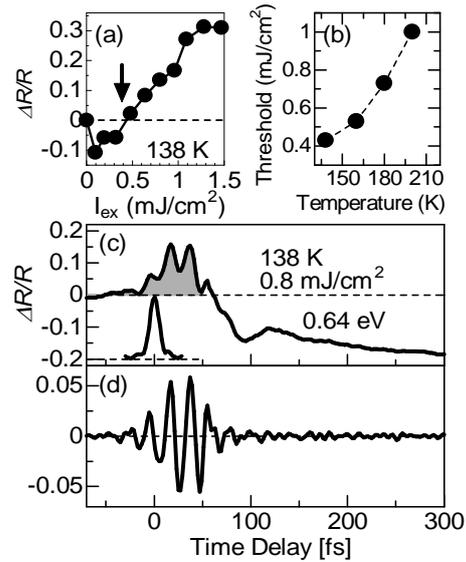


Fig. 3. (a) Excitation intensity (I_{ex}) dependence of $\Delta R/R$ at 0.64 eV (138 K). The threshold intensity is indicated by the arrows. (b) Temperature dependence of the threshold intensity. (c) Time evolution of $\Delta R/R$ at 138 K (0.8 mJ/cm²). The cross correlation profile of the pump and the probe pulse is also shown. (d) Oscillating component obtained by using Fourier filter (>20 THz).

for $I_{\text{ex}} < 0.5 \text{ mJ/cm}^2$ is attributable to the increase of electron temperature. The threshold intensity (shown by arrows in Fig. 3(a)) decreases near the $T_{\text{CO}} = 135 \text{ K}$ as shown in Fig. 3(b). Such threshold behaviour near the CO temperature also supports the interpretation in terms of the photoinduced M-I change.

Figs. 3(c) and 3(d) respectively show the time evolution of $\Delta R/R$ measured at 0.64 eV (Fig. 3(c)) and the high-frequency (> 20 THz) oscillating component (Fig. 3(d)) obtained by Fourier high-pass filter. The oscillation with a period of 20 fs is attributable to the coherent oscillation of the correlated charge [2], because spectral shape obtained by the wavelet (WL) analysis shown in Fig. 1(d) is analogous to that of the inter-molecular charge transfer excitation in optical conductivity (Fig. 1(c)). This coherent charge oscillation is closely related to the M-I change, because the charge oscillation disappears with the decay of the photoinduced insulating state.

Considering that the I-M change can be observed only for $I_{\text{ex}} > 0.5 \text{ mJ/cm}^2$ ($\sim 6 \text{ MV/cm}$), the M-I change is induced by the strong electric field. According to the "dynamical localization" theory [3-5] under strong field, intersite transfer integral (t) is modulated and can be reduced to zero for $\sim 100 \text{ MV/cm}$ [5]. The M-I change is considered to occur as a precursory phenomena of the dynamical localization. Small change of t can result in the M-I change, because t and U , V are competing with each other near the M-I phase boundary in α -(ET)₂I₃.

4. Summary

Photoinduced insulating were realized in organic conductor α -(ET)₂I₃ by strong electric field (10 MV/cm) of 1.5-cycle (7 fs) near infrared pulse. This phenomenon can be understood as a precursory phenomena for the dynamical localization effect near the M-I transition temperature.

5. References

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